



On the well-posedness and the numerical approximation of some problems arising from the modelisation of metamaterials

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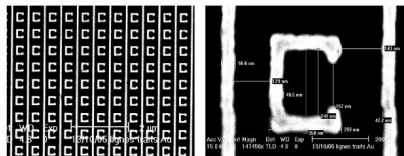
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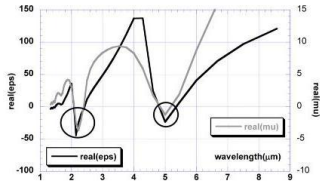
r e t o u r s u r i n n o v a t i o n

Framework :

- Artificial medium created by **homogenization** of small **physical** components.
- **Exotic** behavior at some **frequencies**.



(a) Periodic array of SRR



(b) Real part of homogenized parameters

Fig.: Real part of ϵ and μ of a periodic array of SRR, Kanté B., A. De Lustrac et als, **Metamaterials for optical and radio communications.**

Some application of metamaterials :

- Super-Lens (with materials with **negative refractive index**).
- Alice's mirror,
- Control of light (with **photonic crystals**).
- Invisibility, cloaking
- Seismic protection for buildings.
- ...

Setting of the problem

Hypothesis :

- $\Omega \subset \mathbb{R}^3$ = a **bounded** open set of \mathbb{R}^3 with C^2 boundary with outward unitary normal denoted by \mathbf{n} .
- Metamaterials = **compact** material.

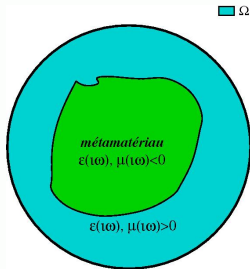


Fig.: Géometric setting of the problem

Maxwell's system in Laplace transform :

$$\left\{ \begin{array}{l} \text{Find } (E, H) \in \mathcal{D}(\mathbb{M}) \text{ such that :} \\ (K(p, x) + \mathbb{M}) \begin{pmatrix} E \\ H \end{pmatrix} = f, \quad L^2(\Omega)^6 \end{array} \right.$$

where :

$$\mathcal{D}(\mathbb{M}) = \left\{ (E, H) \in H(\text{curl}, \Omega)^2 \mid \mathbf{n}(x) \times (E + \Lambda(x)(\mathbf{n}(x) \times H)) = 0, \quad H^{-1/2}(\partial\Omega) \right\}$$

$$p = iw + \eta, \quad \mathbb{M} = \begin{pmatrix} 0 & -\nabla \times \\ \nabla \times & 0 \end{pmatrix}$$

and :

$$K(p, x) = \begin{pmatrix} p\varepsilon(p, x) & 0 \times \mathbb{I}_3 \\ 0 \times \mathbb{I}_3 & p\mu(p, x) \end{pmatrix} \in \text{Hom}(\mathbb{C}^6),$$

$$\forall u \in \mathbb{C}^3, \quad \text{Re} \langle (\Lambda + \Lambda^*)u, \bar{u} \rangle \geq \alpha|u|^2.$$

Main problem

Physical parameters of metamaterials **depends on the frequency** and could become **negative definite** on the behavior of some p .

⇒ The multiplicative operator $K(p)$ is **no longer coercive** for some p .

Goal

- a) Give **conditions on the homogenized parameters** of the metamaterials leading to the **well-posedness** of the Maxwell's system.
- b) Look for numerical scheme suitable for metamaterials.

An example : periodic array of SRR

B. Kanté, SN Burokur, F. Gadot, and A. de Lustrac. Métamatériau à indice de réfraction négatif en infrarouge.

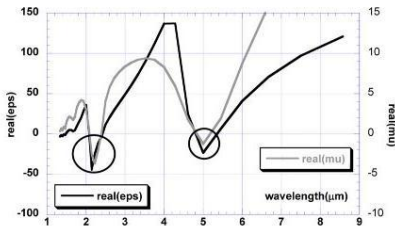


Fig.: Periodic array of S.R.R

Homogenized parameters

$$[\varepsilon(p, x)] = \left(1 + \frac{\omega_p^2}{p^2}\right) \mathbb{I}_3, \quad [\mu(p, x)] = \left(1 + \frac{\delta p^2}{-p^2 - \omega_0^2 + p\Gamma}\right) \mathbb{I}_3.$$

Main result

Assumptions on the metamaterial :

- (**p-Regularity**) $p \in D_0 \longrightarrow K(p, x)$ is holomorphic for almost all $x \in \Omega$.
- (**x-regularity**) $x \in \bar{\Omega} \longrightarrow K(p, x)$ belong to $L^\infty(\bar{\Omega})$ for all $p \in D_0$.
- (**Classical**) There exist $p_0 \in D_0$ such that $\operatorname{Re}(K(p_0) + K(p_0)^*) \succeq \alpha$.
- (**technical...**) There exist $a \in \operatorname{Lip}(\bar{\Omega}, \mathcal{O}(D_0))$ such that $\operatorname{Re}(a(p)K(p) + a(p)^*K(p)^*) \succeq \alpha$.

Theorem

Assume that (**p-Regularity**)-(**x-Regularity**)-(**Classical**)-(**technical...**) are satisfied. Then the Maxwell's system in presence of metamaterial is **well-posed** for all $p \in D_0 \setminus S$ where S is a discrete set of D_0 . Moreover, one has the estimate :

$$\left\{ \int_{\Omega} |E(p, x)|^2 + |H(p, x)|^2 dx \right\}^{\frac{1}{2}} \lesssim \left\{ \int_{\Omega} \left| (K(p, x) + \mathbb{M}) \begin{pmatrix} E(p, x) \\ H(p, x) \end{pmatrix} \right|^2 dx \right\}^{\frac{1}{2}}$$

Numerical simulation

Idea : From [A.K. Aziz, S. Leventhal, Finite element approximation for first order systems, SIAM J Num An, 1978.] \Rightarrow Limited to **coercive** operators.

Variational formulation :

$$\text{Find } \mathbf{u}_h = (\mathbf{E}_h, \mathbf{H}_h) \in \mathcal{V}_h \subset \mathcal{D}(\mathbb{M}) \text{ such that } \forall \psi_h \in (K(p, x) + \mathbb{M}) \mathcal{V}_h : \\ \int_{\Omega} \langle (K(p, x) + \mathbb{M}) \mathbf{u}_h, \overline{\psi}_h \rangle dx = \int_{\Omega} \langle f, \overline{\psi}_h \rangle dx$$

Theorem

Assume that $\forall u \in H^s(\Omega) \cap \mathcal{D}(\mathbb{M})$ there exists $\tilde{u}_h \in \mathcal{V}_h$ such that

$$\| u - \tilde{u}_h \|_{\mathcal{D}(\mathbb{M})} \leq Ch^s \| u \|_{H^s(\Omega) \cap \mathcal{D}(\mathbb{M})} .$$

Then if the solution $u \in H^s(\Omega) \cap \mathcal{D}(\mathbb{M})$ one has the error inequality :

$$\| u - u_h \|_{L^2(\Omega)^6} \leq Ch^s \| u \|_{H^s(\Omega) \cap \mathcal{D}(\mathbb{M})} .$$

- Let $\mathcal{A}(u, v) = \int_{\Omega} \langle (K(\rho, x) + \mathbb{M}) u, \overline{(K(\rho, x) + \mathbb{M}) v} \rangle dx$. From the **theorem**, one gets :

$$\forall u \in \mathcal{D}(\mathbb{M}), \mathcal{A}(u, u) \geq C \|u\|_{L^2(\Omega)}^2 .$$

- Let u_h satisfying $\mathcal{A}(u_h, \phi_h) = \int_{\Omega} \langle f, \overline{\phi_h} \rangle dx$. Since $\dim(\mathcal{V}_h) < \infty$ and $\mathcal{V}_h \subset \mathcal{D}(\mathbb{M})$, u_h is unique.
- Let $\tilde{u}_h \in \mathcal{V}_h$ such that

$$\|u - \tilde{u}_h\|_{\mathcal{D}(\mathbb{M})} \leq Ch^s \|u\|_{H^s(\Omega) \cap \mathcal{D}(\mathbb{M})} .$$

- For all $v_h \in \mathcal{V}_h$ we have :

$$\mathcal{A}(u - \tilde{u}_h, v_h) = \mathcal{A}(u_h - \tilde{u}_h, v_h).$$

- Taking $v_h = u_h - \tilde{u}_h$ and using the **Schwartz inequality**, one gets :

$$\begin{aligned} \| (K(p, x) + \mathbb{M})(u_h - \tilde{u}_h) \|_{L^2(\Omega)} &\leq C \| (K(p, x) + \mathbb{M})(u - \tilde{u}_h) \|_{L^2(\Omega)} \\ &\leq C \| u - \tilde{u}_h \|_{\mathcal{D}(\mathbb{M})} \\ &\leq Ch^s \| u \|_{H^s(\Omega) \cap \mathcal{D}(\mathbb{M})} \end{aligned}$$

- Then, from the **coercivity inequality**, it follows :

$$\| u - u_h \|_{L^2(\Omega)} \leq Ch^s \| u \|_{H^s(\Omega) \cap \mathcal{D}(\mathbb{M})}$$

- $\Omega =]a, b[\subset \mathbb{R}$.

$$\left\{ \begin{array}{l} \text{Find } (u, v) \in H_0^1(\Omega) \times H^1(\Omega) \text{ such that :} \\ \rho \varepsilon(\rho, x) u - \partial_x v = f_1, \\ \rho \mu(\rho, x) v - \partial_x u = f_2 \end{array} \right. \quad (1)$$

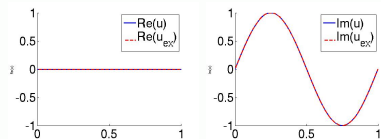
Remark :

Problem **well-posed** if (**x-p-regularity**) and ("**Classical**") hold.

An exact solution :

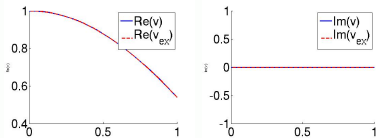
- $\Omega =]0, 1[$, $f_1(x) = i p \varepsilon \sin(2\pi x) + \sin(x)$, $f_2(x) = \rho \mu \cos(x) - 2i\pi \cos(2\pi x)$.
- $u(x) = i \sin(2\pi x)$, $v(x) = \cos(x)$.
- $p = iw$, $w = 15$.

Numerical experiments (1) : the vaccum



(a) $\text{Re}(u)$

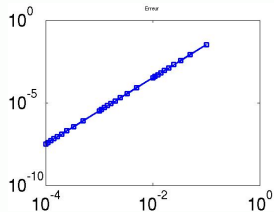
(b) $\text{Im}(u)$



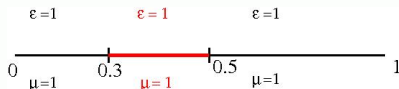
(c) $\text{Re}(v)$

(d) $\text{Im}(v)$

Fig.: The vacuum

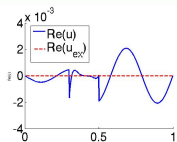


(a) Error, $\|u - u_h\|_{L^2}$

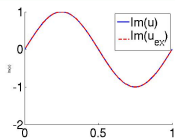


(b) Geometrie

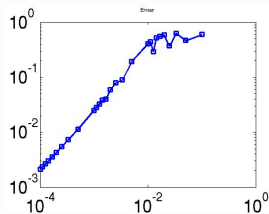
Numerical experiments (2) : A "classical" media



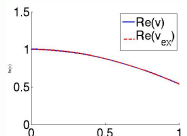
(c) $\text{Re}(u)$



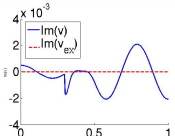
(d) $\text{Im}(u)$



(a) Error, $\|u - u_h\|_{L^2}$

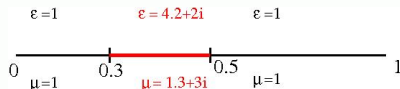


(e) $\text{Re}(v)$



(f) $\text{Im}(v)$

Fig.: A dielectric media



(b) Geometrie

Numerical experiments (3) : A metamaterial

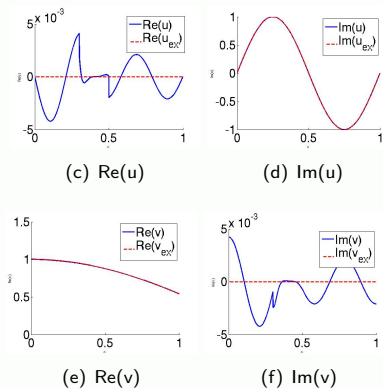
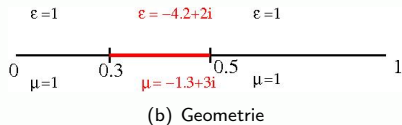
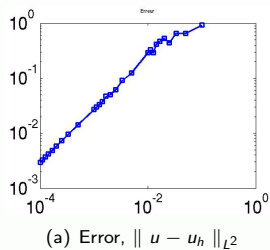


Fig.: A metamaterial



Conclusions and prospects

Conclusion :

- **smooth** metamaterials \implies **well-posed** problems.
- **Numerical** approximation \implies OK with some kind of **finite element method**.

Prospects :

- Study of the **well-posedness** for L^∞ metamaterials (transmission problems ???).
- Convergence of **Discontinuous Galerkin** methods.
- Numerical test in 2D.
-

Thank you for you attention.